# A combinatorial formula for Macdonald polynomials

Arun Ram
Department of Mathematics and Statistics
University of Melbourne
Parkville VIC 3010 Australia
aram@unimelb.edu.au
and
Department of Mathematics
University of Wisconsin
Madison, WI 53706

Martha Yip
Department of Mathematics
University of Wisconsin
Madison, WI 53706
yip@math.wisc.edu

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#### Abstract

Abstract. In this paper we use the combinatorics of alcove walks to give uniform combinatorial formulas for Macdonald polynomials for all Lie types. These formulas resemble the formulas of Haglund-Haiman-Loehr for Macdonald polynomials of type  $GL_n$ . At q=0 these formulas specialize to the formula of Schwer for the Macdonald spherical function in terms of positively folded alcove walks and at q=t=0 these formulas specialize to the formula for the Weyl character in terms of the Littelmann path model (in the positively folded gallery form of Gaussent-Littelmann).

### 1 Introduction

The Macdonald polynomials were introduced in the mid 1980s [Mac1] [Mac2] as a remarkable family of orthogonal polynomials generalizing the spherical functions for a p-adic group, the Weyl characters, the Jack polynomials and the zonal polynomials. In the early 1990s Cherednik [Ch] introduced the double affine Hecke algebra (the DAHA) and used it as a tool to prove conjectures of Macdonald. The DAHA is a fundamental tool for studying Macdonald polynomials. Using the DAHA, the nonsymmetric Macdonald polynomials  $E_{\mu}$  can be constructed by applying products of "intertwining operators"  $\tau_i^{\vee}$  to the generator 1 of the polynomial representation of the DAHA (see [Hai, Prop. 6.13]), and the symmetric Macdonald polynomials  $P_{\mu}$  can then be constructed from the  $E_{\mu}$  by "symmetrizing" (see [Mac3, Remarks after (6.8)]).

Of recent note in the theory of Macdonald polynomials has been the success of Haglund-Haiman-Loehr in giving, in the type  $GL_n$  case, explicit combinatorial formulas for the expansion of Macdonald polynomials in terms of monomials. These formulas were conjectured by J. Haglund and proved by Haglund-Haiman-Loehr in [HHL1] and [HHL2]. The papers [GR] and [Hai] are excellent survey articles discussing these developments.

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Following a key idea of C. Schwer [Sc], the paper [Ra] developed a combinatorics for working in the affine Hecke algebra, the *alcove walk model*. It turns out that this combinatorics is the ideal tool for expansion of products of intertwining operators in the DAHA. These expansions, when applied to the generator of the polynomial representation of the DAHA, give formulas for the Macdonald polynomials for all root systems.

Our formulas resemble the formulas obtained by Haglund-Haiman-Loehr [HHL1] [HHL2] in type  $GL_n$ . The form of the terms is the same but our formula has many more terms. Lenart [Le] has explained how to combine terms in our formula to obtain a formula similar to the Haglund-Haiman-Loehr (HHL) one. The individual terms of our formula have, in general, many more factors appearing in the denominators than the terms in the HHL formula. In the  $GL_n$  case there is an integral form of the Macdonald polynomials obtained by multiplying by a factor. This factor cancels all denominators in the HHL formula but not in ours and so the "integrality" of the integral form is not directly visible from our formula in the way that it is visible in the HHL formula. To our knowledge, a suitable definition of an integral form of the Macdonald polynomials is still unknown in the general root system setting.

At q=0 the symmetric Macdonald polynomials are the Hall-Littlewood polynomials or the Macdonald spherical functions. These are the spherical functions for G/K, where G is a p-adic group and K is a maximal compact subgroup. The work of Schwer [Sc, Thm. 1.1] provided fomulas for the expansion of the Macdonald spherical functions in terms of positively folded alcove walks. See [Ra, Thm. 4.2(a)] for a description of the Schwer formula in terms of the alcove walk model. The formula for Macdonald polynomials which we give in Theorem 3.4 reduces to the Schwer formula at q=0.

At q=t=0 the symmetric Macdonald polynomials are the Weyl characters or Schur functions. In this case our formula for the Macdonald polynomial specializes to the formula for the Weyl character in terms of the Littelmann path model (in the maximal dimensional positively folded gallery form of Gaussent-Littelmann [GL, Cor. 1 p. 62], or the  $\lambda$ -chain form of Lenart-Postnikov [LP1] [LP2]).

It is interesting to note that, in the formulas for the symmetric Macdonald polynomials, the negative folds and the positive folds play an equal role. It is known [GL] that the alcove walks with only positive folds contain detailed information about the geometry of Mirković-Vilonen intersections in the loop Grassmannian. It is tantalizing to wonder whether the alcove walks with both positive and negative folds play a similar role in the geometry of flag varieties for reductive groups over two dimensional local fields and whether the expansions of Macdonald polynomials in this paper are shadows of geometric decompositions.

The papers [GL], [LP1], [LP2], and [Ra] explain how the combinatorics of alcove walks is almost equivalent to the combinatorics of crystal bases and Kashiwara operators (at least for the positively folded alcove walks of maximal dimension). Our expansions of Macdonald polynomials in terms of alcove walks give insight into possible relationships between Macdonald polynomials and crystal and canonical bases.

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## 2 Double Weyl groups, braid groups and Hecke algebras

In this section we review the basic definitions and notations for affine Weyl groups and double affine Hecke algebras following the expositions in [Ra], [Ch], [Mac4] and [Hai]. Following the definitions we prove Theorem 2.2, a formula for the expansion of products of intertwining operators in the DAHA. This formula is a "lift into the DAHA" of the expansions of Macdonald polynomials given in Section 3.

### 2.1 Double affine Weyl groups

Let  $\mathfrak{h}_{\mathbb{Z}}$  be a  $\mathbb{Z}$ -lattice with an action of a finite subgroup  $W_0$  of  $GL(\mathfrak{h}_{\mathbb{Z}})$  generated by reflections. Then  $W_0$  acts on  $\mathfrak{h}_{\mathbb{Z}}^*$  by

$$\langle w\mu, \lambda^{\vee} \rangle = \langle \mu, w^{-1}\lambda^{\vee} \rangle, \quad \text{where} \quad \langle \lambda^{\vee}, \mu \rangle = \mu(\lambda^{\vee}) \quad \text{for } \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}, \, \mu \in \mathfrak{h}_{\mathbb{Z}}^*.$$
 (2.1)

Let  $R^+ \subseteq \mathfrak{h}_{\mathbb{Z}}^*$  and  $(R^{\vee})^+ \subseteq \mathfrak{h}_{\mathbb{Z}}$  denote fixed choices of the positive roots and the positive coroots so that the reflections  $s_{\alpha}$  in  $W_0$  act on  $\mathfrak{h}_{\mathbb{Z}}$  and on  $\mathfrak{h}_{\mathbb{Z}}^*$  by

$$s_{\alpha}\lambda = \lambda - \langle \lambda, \alpha^{\vee} \rangle \alpha$$
 and  $s_{\alpha}\lambda^{\vee} = \lambda^{\vee} - \langle \lambda^{\vee}, \alpha \rangle \alpha^{\vee}$ , respectively. (2.2)

The groups

$$X = \{ X^{\mu} \mid \mu \in \mathfrak{h}_{\mathbb{Z}}^* \} \quad \text{and} \quad Y = \{ Y^{\lambda^{\vee}} \mid \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}} \}$$
 (2.3)

with

$$X^{\mu}X^{\nu} = X^{\mu+\nu}$$
 and  $Y^{\lambda^{\vee}}Y^{\sigma^{\vee}} = Y^{\lambda^{\vee}+\sigma^{\vee}}$  (2.4)

are the groups  $\mathfrak{h}_{\mathbb{Z}}^*$  and  $\mathfrak{h}_{\mathbb{Z}}$  respectively, except written multiplicatively, and the semidirect product

$$W_0 \ltimes (X \times Y) = \{ X^{\mu} w Y^{\lambda^{\vee}} \mid w \in W_0, \mu \in \mathfrak{h}_{\mathbb{Z}}^*, \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}} \}$$
 (2.5)

has additional relations

$$wX^{\mu} = X^{w\mu}w \quad \text{and} \quad wY^{\lambda^{\vee}} = Y^{w\lambda^{\vee}}w,$$
 (2.6)

for  $w \in W_0$ ,  $\mu \in \mathfrak{h}_{\mathbb{Z}}^*$  and  $\lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}$ .

Assume that the action of  $W_0$  on  $\mathfrak{h}_{\mathbb{C}} = \mathbb{C} \otimes_{\mathbb{Z}} \mathfrak{h}_{\mathbb{Z}}$  is irreducible. The double affine Weyl group  $\widetilde{W}$  is the universal central extension of  $W_0 \ltimes (X \times Y)$ . If e is the smallest integer such that  $\langle \lambda^{\vee}, \mu \rangle \in \frac{1}{e}\mathbb{Z}$  for all  $\lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}$  and  $\mu \in \mathfrak{h}_{\mathbb{Z}}^*$  then  $\widetilde{W}$  is presented by

$$\widetilde{W} = \{q^k X^{\mu} w Y^{\lambda^{\vee}} \mid k \in \frac{1}{e} \mathbb{Z}, \mu \in \mathfrak{h}_{\mathbb{Z}}^*, \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}, w \in W_0\}$$

with (2.4), (2.6) and

$$X^{\mu}Y^{\lambda^{\vee}} = q^{\langle \lambda^{\vee}, \mu \rangle} Y^{\lambda^{\vee}} X^{\mu}, \quad \text{for } \mu \in \mathfrak{h}_{\mathbb{Z}}^{*}, \ \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}.$$
 (2.7)

The subgroup  $\{q^k X^{\mu} Y^{\lambda^{\vee}} \mid k \in \frac{1}{e} \mathbb{Z}, \mu \in \mathfrak{h}_{\mathbb{Z}}^*, \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}} \}$  is a Heisenberg group and

$$W = \{X^{\mu}w \mid \mu \in \mathfrak{h}_{\mathbb{Z}}^*, w \in W_0\} \quad \text{and} \quad W^{\vee} = \{wY^{\lambda^{\vee}} \mid \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}, w \in W_0\} \quad (2.8)$$

are affine Weyl groups inside  $\widetilde{W}$ . Letting

$$q = X^{\delta} = Y^{-d} \tag{2.9}$$

and extending the notation of (2.6) gives actions of  $W^{\vee}$  on  $\mathfrak{h}_{\mathbb{Z}}^* \oplus \mathbb{Z}\delta$  and W on  $\mathfrak{h}_{\mathbb{Z}} \oplus \mathbb{Z}d$  with

$$Y^{\lambda^{\vee}}\mu = \mu - \langle \mu, \lambda^{\vee} \rangle \delta$$
 and  $X^{\mu}\lambda^{\vee} = \lambda^{\vee} - \langle \lambda^{\vee}, \mu \rangle d$ . (2.10)

Let  $\varphi \in R$  be the highest root and  $\varphi^{\vee} \in R^{\vee}$  the highest coroot and let

$$s_0 = Y^{\varphi^{\vee}} s_{\varphi} \quad \text{and} \quad s_0^{\vee} = X^{\varphi} s_{\varphi^{\vee}}.$$
 (2.11)

Let

$$\alpha_0 = -\varphi + \delta, \quad \alpha_0^{\vee} = -\varphi^{\vee} + d, \qquad \langle d, \mu \rangle = 0, \quad \langle \lambda^{\vee}, \delta \rangle = 0, \qquad \langle d, \delta \rangle = 0,$$
 (2.12)

so that

$$s_0 \mu = \mu - \langle \mu, \alpha_0^{\vee} \rangle \alpha_0$$
 and  $s_0^{\vee} \lambda^{\vee} = \lambda^{\vee} - \langle \lambda^{\vee}, \alpha_0 \rangle \alpha_0^{\vee}$ . (2.13)

The alcoves of  $\mathfrak{h}_{\mathbb{R}}^*=\mathbb{R}\otimes_{\mathbb{Z}}\mathfrak{h}_{\mathbb{Z}}^*$  are the connected components of

$$\mathfrak{h}_{\mathbb{R}}^* \setminus \left( \bigcup_{\alpha^{\vee} \in (R^{\vee})^+, j \in \mathbb{Z}} \mathfrak{h}^{\alpha^{\vee} + jd} \right) \quad \text{where} \quad \mathfrak{h}^{\alpha^{\vee} + jd} = \{ x \in \mathfrak{h}_{\mathbb{R}}^* \mid \langle x, \alpha^{\vee} \rangle = -j \}.$$
 (2.14)

The action of  $W=\{X^{\mu}w\mid \mu\in\mathfrak{h}_{\mathbb{Z}}^*, w\in W_0\}$  on  $\mathfrak{h}_{\mathbb{R}}^*$  given by

$$X^{\mu} \cdot \nu = \nu + \mu$$
 and  $w \cdot \nu = w\nu$ , for  $w \in W_0$ ,  $\mu \in \mathfrak{h}_{\mathbb{Z}}^*$  and  $\nu \in \mathfrak{h}_{\mathbb{R}}^*$ , (2.15)

sends alcoves to alcoves;  $s_0^{\vee}, \dots, s_n^{\vee}$  are the reflections in the walls  $\mathfrak{h}^{\alpha_0^{\vee}}, \dots, \mathfrak{h}^{\alpha_n^{\vee}}$  of the fundamental alcove

$$1 = \{ x \in \mathfrak{h}_{\mathbb{R}}^* \mid \langle x, \alpha_i^{\vee} \rangle \ge 0, \text{ for } i = 0, 1, \dots, n \}; \quad \text{and}$$
 (2.16)

$$\ell(v) = \text{(number of hyperplanes between 1 and } v)$$
 (2.17)

is the *length* of  $v \in W$ . Let  $\Omega^{\vee}$  be the set of length zero elements of W. The affine Weyl group W has an alternate presentation by generators  $s_0^{\vee}, s_1^{\vee}, \ldots, s_n^{\vee}$  and  $\Omega^{\vee}$  with relations

$$(s_i^{\vee})^2 = 1, \qquad \underbrace{s_i^{\vee} s_j^{\vee} \cdots}_{m_{ij}^{\vee}} = \underbrace{s_j^{\vee} s_i^{\vee} \cdots}_{m_{ij}^{\vee}}, \quad \text{and} \quad g^{\vee} s_i^{\vee} (g^{\vee})^{-1} = s_{\sigma^{\vee}(i)}^{\vee}, \quad \text{for } g^{\vee} \in \Omega^{\vee}, \quad (2.18)^{\vee}$$

where  $\pi/m_{ij}^{\vee}$  is the angle between  $\mathfrak{h}^{\alpha_i^{\vee}}$  and  $\mathfrak{h}^{\alpha_j^{\vee}}$  and  $\sigma^{\vee}$  denotes the permutation of the  $\mathfrak{h}^{\alpha_i^{\vee}}$  induced by the action of  $g^{\vee}$ . If  $\Omega^{\vee} \times \mathfrak{h}_{\mathbb{R}}^*$  is  $|\Omega^{\vee}|$  copies of  $\mathfrak{h}_{\mathbb{R}}^*$  (sheets), with  $\Omega^{\vee}$  acting by switching sheets then there is a bijection

$$W \longleftrightarrow \{\text{alcoves in } \Omega^{\vee} \times \mathfrak{h}_{\mathbb{R}}^*\}$$
 (2.19)

and we will often identify  $v \in W$  with the corresponding alcove in  $\Omega^{\vee} \times \mathfrak{h}_{\mathbb{R}}^*$ . The pictures illustrating this bijection in type  $SL_3$  are displayed in the appendix.

The periodic orientation is the orientation of the hyperplanes  $\mathfrak{h}^{\alpha^{\vee}+kd}$  such that

(a) 1 is on the positive side of 
$$\mathfrak{h}^{\alpha^{\vee}}$$
 for  $\alpha^{\vee} \in (R^{\vee})^+$ , (2.20)

(b)  $\mathfrak{h}^{\alpha^{\vee}+kd}$  and  $\mathfrak{h}^{\alpha^{\vee}}$  have parallel orientations.

The pictures in the appendix illustrate the periodic orientation for type  $SL_3$ .

A similar "pictorial" viewpoint applies to the group  $W^{\vee}$  acting on  $\Omega \times \mathfrak{h}_{\mathbb{R}}$  where  $\mathfrak{h}_{\mathbb{R}} = \mathbb{R} \otimes_{\mathbb{Z}} \mathfrak{h}_{\mathbb{Z}}$  and  $\Omega$  is the set of length zero elements of  $W^{\vee}$ . Then  $W^{\vee}$  has an alternate presentation by generators  $s_0, s_1, \ldots, s_n$  and  $\Omega$  with relations

$$s_i^2 = 1,$$
  $\underbrace{s_i s_j \cdots}_{m_{ij}} = \underbrace{s_j s_i \cdots}_{m_{ij}},$  and  $g s_i g^{-1} = s_{\sigma(i)},$  for  $g \in \Omega,$  (2.21)

where  $\pi/m_{ij}$  is the angle between  $\mathfrak{h}^{\alpha_i}$  and  $\mathfrak{h}^{\alpha_j}$  and  $\sigma$  denotes the permutation of the  $\mathfrak{h}^{\alpha_i}$  induced by the action of g.

### 2.2 Double affine braid groups

The double affine braid group  $\tilde{\mathcal{B}}$  is the group generated by  $T_0, \ldots, T_n, \Omega$  and X with relations

$$\underbrace{T_i T_j \cdots}_{m_{ij}} = \underbrace{T_j T_i \cdots}_{m_{ij}}, \quad g T_i g^{-1} = T_{\sigma(i)}, \quad g X^{\mu} = X^{g\mu} g, \tag{2.22}$$

for  $g \in \Omega$ , and

$$T_i X^{\mu} = X^{s_i \mu} T_i, \quad \text{if } \langle \mu, \alpha_i^{\vee} \rangle = 0, T_i X^{\mu} T_i = X^{s_i \mu}, \quad \text{if } \langle \mu, \alpha_i^{\vee} \rangle = 1,$$
 for  $i = 0, 1, \dots, n,$  (2.23)

where the action of  $W^{\vee}$  on  $\mathfrak{h}_{\mathbb{Z}}^* \oplus \mathbb{Z}\delta$  is as in (2.10). The element

$$q = X^{\delta}$$
 is in the center of  $\tilde{\mathcal{B}}$ . (2.24)

For  $w \in W^{\vee}$ , view a reduced word  $w = gs_{i_1} \cdots s_{i_{\ell}}$  as a minimal length path p from the fundamental alcove to w in  $\mathfrak{h}_{\mathbb{R}}$  and define

$$Y^{w} = g(T_{i_1})^{\epsilon_1} \cdots (T_{i_\ell})^{\epsilon_\ell}, \quad \text{with} \quad \epsilon_k = \begin{cases} +1, & \text{if the } k \text{th step of } p \text{ is} & \frac{-}{} & + \\ -1, & \text{if the } k \text{th step of } p \text{ is} & \frac{-}{} & + \end{cases}, \quad (2.25)$$

with respect to the *periodic orientation* (see (2.20) and the pictures in the appendix). For  $v \in W$ , view a reduced word  $v = g^{\vee} s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  as a minimal length path  $p^{\vee}$  from the fundamental alcove to v in  $\mathfrak{h}_{\mathbb{R}}^*$  and define

$$X^{v} = g^{\vee}(T_{i_{1}}^{\vee})^{\epsilon_{1}^{\vee}} \cdots (T_{i_{\ell}}^{\vee})^{\epsilon_{\ell}^{\vee}}, \quad \text{with} \quad \epsilon_{k}^{\vee} = \begin{cases} -1, & \text{if the } k \text{th step of } p^{\vee} \text{ is} & \xrightarrow{-} & + \\ +1, & \text{if the } k \text{th step of } p^{\vee} \text{ is} & \xrightarrow{-} & + \end{cases}, \tag{2.26}$$

Let  $T_i^{\vee} = T_i$ , for i = 1, 2, ..., n,

$$g^{\vee} = X^{\omega_g} T_{w_q w_0}^{\vee}, \qquad (T_0^{\vee})^{-1} = X^{\varphi} T_{s_{\varphi}}^{\vee}, \qquad g = Y^{\omega_g^{\vee}} T_{w_0 w_q}^{-1}, \qquad T_0 = Y^{\varphi^{\vee}} T_{s_{\varphi}}^{-1}. \tag{2.27}$$

where  $\varphi$  and  $\varphi^{\vee}$  are as in (2.11) and, using the action in (2.15),  $\omega_g = g^{\vee} \cdot 0$  and  $w_g$  is the longest element of the stabilizer of  $\omega_g$  in  $W_0$ .

The following theorem, discovered by Cherednik [Ch, Thm. 2.2], is proved in [Mac4, 3.5-3.7], in [Io], and in [Hai, 4.13-4.18].

**Theorem 2.1.** (Duality) Let  $Y^d = q^{-1}$ . The double affine braid group  $\tilde{\mathcal{B}}$  is generated by  $T_0^{\vee}, T_1^{\vee}, \dots, T_n^{\vee}, \Omega^{\vee}$  and Y with relations

$$\underbrace{T_{j}^{\vee}T_{j}^{\vee}\cdots}_{m_{ij}^{\vee}} = \underbrace{T_{j}^{\vee}T_{i}^{\vee}\cdots}_{m_{ij}^{\vee}}, \qquad g^{\vee}T_{i}^{\vee}(g^{\vee})^{-1} = T_{\sigma^{\vee}(i)}^{\vee}, \qquad g^{\vee}Y^{\lambda^{\vee}} = Y^{g^{\vee}\lambda^{\vee}}g^{\vee}, \tag{2.28}$$

for  $g^{\vee} \in \Omega^{\vee}$ , and

$$T_{i}^{\vee}Y^{\lambda^{\vee}} = Y^{s_{i}^{\vee}\lambda^{\vee}}T_{i}^{\vee}, \quad if \langle \lambda^{\vee}, \alpha_{i} \rangle = 0, (T_{i}^{\vee})^{-1}Y^{\lambda^{\vee}}(T_{i}^{\vee})^{-1} = Y^{s_{i}^{\vee}\lambda^{\vee}}, \quad if \langle \lambda^{\vee}, \alpha_{i} \rangle = 1,$$
 for  $i = 0, 1, \dots, n$ , (2.29)

where the action of W on  $\mathfrak{h}_{\mathbb{Z}} \oplus \mathbb{Z}d$  is as in (2.10).

### 2.3 Double affine Hecke algebras

In the following, for simplicity of exposition we shall assume that we are not in the special case of [Mac4, (2.1.6)] where the root system is type  $C_n$  and  $\mathfrak{h}_{\mathbb{Z}}$  is the (co)root lattice. All our results and proofs are valid in this special case but the definition of the double affine Hecke algebra and the formulas in (2.32) and (2.34) may need some slight modification. See Remark 2.3 for details.

Let  $R^{\vee} = (R^{\vee})^+ \cup (-(R^{\vee})^+)$  be the set of coroots and fix parameters  $c_{\beta^{\vee}}$ , indexed by  $\beta^{\vee} \in R^{\vee} + \mathbb{Z}d$ , such that for all  $w \in W$  and  $\beta^{\vee} \in R^{\vee} + \mathbb{Z}d$ ,

$$c_{\beta^{\vee}} = c_{w\beta^{\vee}}.$$
 Set  $t_{\beta^{\vee}} = q^{c_{\beta^{\vee}}}$  and  $t_i = t_{\alpha_i^{\vee}}.$  (2.30)

The double affine Hecke algebra  $\widetilde{H}$  is the group algebra  $\mathbb{C}\widetilde{\mathcal{B}}$  of the double braid group with the additional relations

$$T_i^2 = (t_i^{1/2} - t_i^{-1/2})T_i + 1, \quad \text{for } i = 0, 1, \dots, n.$$
 (2.31)

The double affine Hecke algebra  $\widetilde{H}$  has bases

$$\{T_wX^\mu\mid w\in W,\ \mu\in\mathfrak{h}_\mathbb{Z}^*\oplus\mathbb{Z}\delta\},\qquad \{Y^{\lambda^\vee}T_w^\vee\mid w\in W^\vee,\ \lambda^\vee\in\mathfrak{h}_\mathbb{Z}\oplus\mathbb{Z}d\},$$

and

$$\{q^k X^{\mu} T_w Y^{\lambda^{\vee}} \mid w \in W_0, \ \lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}, \mu \in \mathfrak{h}_{\mathbb{Z}}^*, k \in \frac{1}{e} \mathbb{Z}\}$$

(see [Hai, Prop. 5.4 and Cor. 5.8]).

In the presence of (2.31) the relations (2.29) are equivalent to

$$T_{i}^{\vee}Y^{\lambda^{\vee}} = Y^{s_{i}^{\vee}\lambda^{\vee}}T_{i}^{\vee} + (t_{i}^{\frac{1}{2}} - t_{i}^{-\frac{1}{2}}) \frac{Y^{\lambda^{\vee}} - Y^{s_{i}^{\vee}\lambda^{\vee}}}{1 - Y^{-\alpha_{i}^{\vee}}}, \quad \text{for } i = 0, 1, \dots, n.$$
 (2.32)

In turn (2.32) is equivalent to

$$\tau_i^{\vee} Y^{\lambda^{\vee}} = Y^{s_i^{\vee} \lambda^{\vee}} \tau_i^{\vee}, \quad \text{for } i = 0, 1, \dots, n,$$
 (2.33)

where

$$\tau_i^{\vee} = T_i^{\vee} + \frac{t_i^{-\frac{1}{2}}(1 - t_i)}{1 - Y^{-\alpha_i^{\vee}}} = (T_i^{\vee})^{-1} + \frac{t_i^{-\frac{1}{2}}(1 - t_i)Y^{-\alpha_i^{\vee}}}{1 - Y^{-\alpha_i^{\vee}}}.$$
 (2.34)

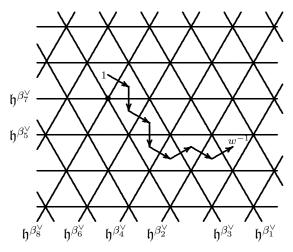
Using that the  $\tau_i^{\vee}$  satisfy the braid relations and that

$$g^\vee Y^{\lambda^\vee} = Y^{g^\vee \lambda^\vee} g^\vee, \qquad \text{write} \qquad \tau_w^\vee Y^{\lambda^\vee} = Y^{w\lambda^\vee} \tau_w^\vee, \quad \text{for } w \in W.$$

Let  $w \in W$  and let  $w = s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  be a reduced word for w. For  $k = 1, \dots, \ell$  let

$$\beta_k^{\vee} = s_{i_{\ell}}^{\vee} s_{i_{\ell-1}}^{\vee} \cdots s_{i_{k+1}}^{\vee} \alpha_{i_k}^{\vee} \quad \text{and} \quad t_{\beta_k^{\vee}} = t_{i_k}, \tag{2.35}$$

so that the sequence  $\beta_{\ell}^{\vee}, \beta_{\ell-1}^{\vee}, \dots, \beta_{1}^{\vee}$  is the sequence of labels of the hyperplanes crossed by the walk  $w^{-1} = s_{i_{\ell}}^{\vee} s_{i_{\ell-1}}^{\vee} \cdots s_{i_{1}}^{\vee}$ . For example, in Type  $A_{2}$ , with  $w = s_{2}^{\vee} s_{0}^{\vee} s_{1}^{\vee} s_{2}^{\vee} s_{1}^{\vee} s_{0}^{\vee} s_{2}^{\vee} s_{1}^{\vee}$  the picture is



Let  $v \in W^{\vee}$ . An alcove walk of type  $i_1, \ldots, i_{\ell}$  beginning at v is a sequence of steps, where a step of type j is

positive j-crossing negative j-crossing positive j-fold negative j-fold

Let  $\mathcal{B}(v, \vec{w})$  be the set of alcove walks of type  $\vec{w} = (i_1, \dots, i_\ell)$  beginning at v. For a walk  $p \in \mathcal{B}(v, \vec{w})$  let

$$f^{+}(p) = \{k \mid \text{the } k \text{th step of } p \text{ is a positive fold}\},\$$
  
 $f^{-}(p) = \{k \mid \text{the } k \text{th step of } p \text{ is a negative fold}\},$  (2.36)

and

$$\operatorname{end}(p) = \operatorname{endpoint} \operatorname{of} p \quad \text{(an element of } W\text{)}.$$
 (2.37)

**Theorem 2.2.** Let  $v, w \in W$ , let  $w = s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  be a reduced word for w and let  $\beta_{\ell}^{\vee}, \ldots, \beta_{1}^{\vee}$  be as defined in (2.35). Then, in  $\widetilde{H}$ ,

$$X^{v}\tau_{w}^{\vee} = \sum_{p \in \mathcal{B}(v, \vec{w})} X^{\text{end}(p)} \left( \prod_{k \in f^{+}(p)} \frac{t_{\beta_{k}^{\vee}}^{-1/2} (1 - t_{\beta_{k}^{\vee}})}{1 - Y^{-\beta_{k}^{\vee}}} \right) \left( \prod_{k \in f^{-}(p)} \frac{t_{\beta_{k}^{\vee}}^{-1/2} (1 - t_{\beta_{k}^{\vee}}) Y^{-\beta_{k}^{\vee}}}{1 - Y^{-\beta_{k}^{\vee}}} \right),$$

where the sum is over all alcove walks of type  $\vec{w} = (i_1, \dots, i_\ell)$  beginning at v.

*Proof.* The proof is by induction on the length of w, the base case being the formulas in (2.34). To do the induction step let  $p \in \mathcal{B}(v, \vec{w})$ ,

$$F^{+}(p) = \left(\prod_{k \in f^{+}(p)} \frac{t_{\beta_{k}^{\vee}}^{-1/2} (1 - t_{\beta_{k}^{\vee}})}{1 - Y^{-\beta_{k}^{\vee}}}\right), \qquad F^{-}(p) = \left(\prod_{k \in f^{-}(p)} \frac{t_{\beta_{k}^{\vee}}^{-1/2} (1 - t_{\beta_{k}^{\vee}}) Y^{-\beta_{k}^{\vee}}}{1 - Y^{-\beta_{k}^{\vee}}}\right)$$

and let

 $p_1, p_2 \in \mathcal{B}(v, \vec{w}s_j)$  be the two extensions of p by a step of type j

(by a crossing and a fold, respectively). Let  $z = \operatorname{end}(p)$ . By induction, a term in  $X^v \tau_w^{\vee} \tau_i^{\vee}$  is

$$X^{z}F^{+}(p)F^{-}(p)\tau_{j}^{\vee} = X^{z}\tau_{j}^{\vee}\left(s_{j}F^{+}(p)\right)\left(s_{j}F^{-}(p)\right)$$

$$= \begin{cases} X^{z}\left(T_{j}^{\vee} + \frac{t_{j}^{-1/2}(1-t_{j})}{1-Y^{-\alpha_{j}^{\vee}}}\right)\left(s_{j}F^{+}(p)\right)\left(s_{j}F^{-}(p)\right), & \text{if } X^{zs_{j}} = X^{z}T_{j}^{\vee}, \\ X^{z}\left((T_{j}^{\vee})^{-1} + \frac{t_{j}^{-1/2}(1-t_{j})Y^{-\alpha_{j}^{\vee}}}{1-Y^{-\alpha_{j}^{\vee}}}\right)\left(s_{j}F^{+}(p)\right)\left(s_{j}F^{-}(p)\right), & \text{if } X^{zs_{j}} = X^{z}(T_{j}^{\vee})^{-1}, \\ = X^{\text{end}(p_{1})}F^{+}(p_{1})F^{-}(p_{1}) + X^{\text{end}(p_{2})}F^{+}(p_{2})F^{-}(p_{2}). \end{cases}$$

The last step of  $p_2$  is

$$\stackrel{-}{\underset{zs_j}{\blacktriangleright}} \stackrel{+}{\underset{z}{\blacktriangleright}} \text{ if } X^{zs_j} = X^z T_j^{\vee}, \quad \text{ and } \quad \stackrel{-}{\underset{zs_j}{\blacktriangleright}} \stackrel{+}{\underset{zs_j}{\blacktriangleright}} \text{ if } X^{zs_j} = X^z (T_j^{\vee})^{-1}.$$

Remark 2.3. In some special cases when the affine root system is nonreduced (see [Mac4, (2.1.6)]) the formulas in (2.32) and (2.34) need modification and the definition of the double affine Hecke algebra may need an additional relation. The most involved of these cases is type  $(C_n^{\vee}, C_n)$  (see [Mac4, (1.4.3)]) where the double affine Hecke algebra needs additional parameters  $u_0^{1/2}$  and  $u_n^{1/2}$  (in the notation of [Mac4],  $u_0^{1/2} = \tau_0'$  and  $u_n^{1/2} = \tau_n'$ ) and additional relations

$$(T_0' - u_0^{\frac{1}{2}})(T_0' + u_0^{-\frac{1}{2}}) = 0 \quad \text{and} \quad (T_0^{\vee} - u_n^{\frac{1}{2}})(T_0^{\vee} + u_n^{-\frac{1}{2}}), \quad \text{where } T_0' = q^{-\frac{1}{2}}X^{\varepsilon_1}T_0^{-1}, \quad (2.38)$$

and the formulas for  $\tau_n^\vee$  and  $\tau_0^\vee$  need to be changed to

$$\tau_{n}^{\vee} = T_{n} + \frac{t_{n}^{-\frac{1}{2}}(1 - t_{n}) + t_{0}^{-\frac{1}{2}}(1 - t_{0})Y^{-\varepsilon_{n}^{\vee}}}{1 - Y^{-2\varepsilon_{n}^{\vee}}} = T_{n}^{-1} + \frac{(t_{n}^{-\frac{1}{2}}(1 - t_{n}) + t_{0}^{-\frac{1}{2}}(1 - t_{0})Y^{\varepsilon_{n}^{\vee}})Y^{-2\varepsilon_{n}^{\vee}}}{1 - Y^{-2\varepsilon_{n}^{\vee}}},$$
(2.39)

and

$$\tau_0^{\vee} = T_0^{\vee} + \frac{u_n^{-\frac{1}{2}}(1 - u_n) + u_0^{-\frac{1}{2}}(1 - u_0)q^{\frac{1}{2}}Y\varepsilon_1^{\vee}}{1 - q^{-1}Y^{-2\varepsilon_1^{\vee}}}$$
(2.40)

$$= (T_0^{\vee})^{-1} + \frac{(u_n^{-\frac{1}{2}}(1 - u_n) + u_0^{-\frac{1}{2}}(1 - u_0)q^{-\frac{1}{2}}Y^{-\varepsilon_1^{\vee}})q^{-1}Y^{-2\varepsilon_1^{\vee}}}{1 - q^{-1}Y^{-2\varepsilon_1^{\vee}}}.$$
 (2.41)

The statement and proof of the analogue of Theorem 2.2 for this case is the same, except with the factors associated to the 0-folds and n-folds replaced by the rational functions in Y which appear in the expressions of  $\tau_0^{\vee}$  and  $\tau_n^{\vee}$  in Equations (2.40), (2.41), and (2.39).

# 3 Macdonald polynomials

In this section we use Theorem 2.2 to give expansions of the nonsymmetric Macdonald polynomials  $E_{\mu}$  (Theorem 3.1) and the symmetric Macdonald polynomials  $P_{\mu}$  (Theorem 3.4).

Let  $\widetilde{H}$  be the double affine Hecke algebra (defined in (2.31)) and let H be the subalgebra of  $\widetilde{H}$  generated by  $T_0, \ldots, T_n$  and  $\Omega$ . The polynomial representation of  $\widetilde{H}$  is

$$\mathbb{C}[X] = \operatorname{Ind}_{H}^{\widetilde{H}}(\mathbf{1}) = \mathbb{C}\operatorname{-span}\{q^{k}X^{\mu}\mathbf{1} \mid k \in \frac{1}{e}\mathbb{Z}, \mu \in \mathfrak{h}_{\mathbb{Z}}^{*}\}$$
(3.1)

with

$$T_i \mathbf{1} = t_i^{1/2} \mathbf{1}$$
 and  $g \mathbf{1} = \mathbf{1}$ , for  $g \in \Omega$ . (3.2)

The monomials  $X^{\mu}\mathbf{1}$ ,  $\mu \in \mathfrak{h}_{\mathbb{Z}}^*$ , form a  $\mathbb{C}[q^{\pm 1/e}]$ -basis of  $\mathbb{C}[X]$ . Another favourite  $\mathbb{C}[q^{\pm 1/e}]$ -basis of  $\mathbb{C}[X]$  is the basis of nonsymmetric Macdonald polynomials

$${E_{\mu} \mid \mu \in \mathfrak{h}_{\mathbb{Z}}^*}, \quad \text{where} \quad E_{\mu} = \tau_{X^{\mu}m}^{\vee} \mathbf{1}$$
 (3.3)

with  $X^{\mu}m$  the minimal length element in the coset  $X^{\mu}W_0$ . Note that  $\tau_w^{\vee}\mathbf{1}=0$  for  $w\in W_0$  since  $\tau_i^{\vee}\mathbf{1}=0$  for  $i=1,2,\ldots,n$ .

If  $\mathfrak{h}_{\mathbb{Z}}^+$  is the set of dominant integral coweights (analogous to  $(\mathfrak{h}_{\mathbb{Z}}^*)^+$  defined in (3.8)),  $\lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}^+$  and  $Y^{\lambda^{\vee}} = s_{i_1} \cdots s_{i_{\ell}}$  is a reduced word, then

$$Y^{\lambda^{\vee}} \mathbf{1} = T_{i_1} \cdots T_{i_{\ell}} \mathbf{1} = t_{i_1}^{\frac{1}{2}} \cdots t_{i_{\ell}}^{\frac{1}{2}} \mathbf{1} = q^{\frac{1}{2}(c_{i_1} + \cdots + c_{i_{\ell}})} \mathbf{1} = q^{\frac{1}{2} \sum_{\alpha \in R^+} c_{\alpha} \langle \lambda^{\vee}, \alpha \rangle} \mathbf{1} = q^{\langle \lambda^{\vee}, \rho_c \rangle} \mathbf{1}$$

since  $\langle \lambda^{\vee}, \alpha \rangle$  is the number of hyperplanes parallel to  $\mathfrak{h}^{\alpha}$  which are between  $Y^{\lambda^{\vee}}$  and 1. If  $\lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}$  then  $\lambda^{\vee} = \mu^{\vee} - \nu^{\vee}$  for some  $\mu^{\vee}$ ,  $\nu^{\vee} \in \mathfrak{h}_{\mathbb{Z}}^+$  and so, for all  $\lambda^{\vee} \in \mathfrak{h}_{\mathbb{Z}}$ ,

$$Y^{\lambda^{\vee}} \mathbf{1} = q^{\langle \lambda^{\vee}, \rho_c \rangle} \mathbf{1}, \quad \text{where} \quad \rho_c = \frac{1}{2} \sum_{\alpha \in R^+} c_{\alpha} \alpha.$$
 (3.4)

More generally, if  $X^{\mu}m$  is the minimal length element of the coset  $X^{\mu}W_0$  then

$$Y^{\lambda^{\vee}} E_{\mu} = Y^{\lambda^{\vee}} \tau_{X^{\mu}m}^{\vee} \mathbf{1} = \tau_{X^{\mu}m}^{\vee} Y^{m^{-1}X^{-\mu}\lambda^{\vee}} \mathbf{1} = \tau_{X^{\mu}m}^{\vee} Y^{m^{-1}(\lambda^{\vee} + \langle \lambda^{\vee}, \mu \rangle d)} \mathbf{1}$$

$$= \tau_{X^{\mu}m}^{\vee} Y^{m^{-1}\lambda^{\vee}} q^{-\langle \lambda^{\vee}, \mu \rangle} \mathbf{1} = q^{\langle m^{-1}\lambda^{\vee}, \rho_c \rangle - \langle \lambda^{\vee}, \mu \rangle} \tau_{X^{\mu}m}^{\vee} \mathbf{1} = q^{\langle \lambda^{\vee}, m \rho_c - \mu \rangle} E_{\mu}$$

$$= q^{\langle \lambda^{\vee}, X^{-\mu}m \cdot \rho_c \rangle} E_{\mu},$$

where, in the last line, the action of W on  $\mathfrak{h}_{\mathbb{Z}}^*$  is as in (2.15). Thus the  $E_{\mu}$  are eigenvectors for the action of the  $Y^{\lambda^{\vee}}$  on the polynomial representation  $\mathbb{C}[X]$ .

Retain the notation of (2.36-2.37) so that if  $w = s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  is a reduced word then  $\mathcal{B}(v, \vec{w})$  denotes the set of alcove walks of type  $\vec{w} = (i_1, \dots, i_{\ell})$  beginning at v. For  $p \in \mathcal{B}(v, \vec{w})$  define the weight wt(p) and the final direction  $\varphi(p)$  of p by

$$X^{\operatorname{end}(p)} = X^{\operatorname{wt}(p)} T_{\varphi(p)}^{\vee}, \quad \text{with } \operatorname{wt}(p) \in \mathfrak{h}_{\mathbb{Z}}^{*} \text{ and } \varphi(p) \in W_{0}.$$
 (3.5)

In other words, wt(p) is the "hexagon where p ends". For  $w \in W$  define

$$t_w^{1/2} = t_{i_1}^{1/2} \dots t_{i_\ell}^{1/2}, \quad \text{if } w = s_{i_1}^{\vee} \dots s_{i_\ell}^{\vee} \text{ is a reduced word.}$$
 (3.6)

If  $\beta_k^\vee = s_{i_\ell}^\vee \cdots s_{i_{k+1}}^\vee \alpha_{i_k}^\vee$  are as defined in (2.35) then, by (3.4),

$$Y^{-\beta_k^\vee}\mathbf{1} = Y^{-(-\gamma^\vee + jd)}\mathbf{1} = q^j q^{\langle \gamma^\vee, \rho_c \rangle}\mathbf{1}, \qquad \text{if } \beta_k^\vee = -\gamma^\vee + jd$$

with  $\gamma^{\vee} \in R^{\vee}$ ,  $j \in \mathbb{Z}$ . By (2.30) and the definition of  $\rho_c$  in (3.4), the constant  $q^{\langle \gamma^{\vee}, \rho_c \rangle}$  is a monomial in the symbols  $t_i^{1/2}$ . To simplify the notation for these constants write  $q^j q^{\langle \gamma^{\vee}, \rho_c \rangle} = q^{\langle -\beta_k^{\vee}, \rho_c \rangle}$  so that

$$Y^{-\beta_k^{\vee}} \mathbf{1} = q^{\langle -\beta_k^{\vee}, \rho_c \rangle} \mathbf{1}. \tag{3.7}$$

**Theorem 3.1.** Let  $\mu \in \mathfrak{h}_{\mathbb{Z}}^*$  and let  $w = X^{\mu}m$  be the minimal length element in the coset  $X^{\mu}W_0$ . Fix a reduced word  $\vec{w} = s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  for w and let  $\beta_{\ell}^{\vee}, \ldots, \beta_{1}^{\vee}$  be as defined in (2.35). With notations as in (3.5-3.7) the nonsymmetric Macdonald polynomial

$$E_{\mu} = \sum_{p \in \mathcal{B}(\vec{\mu})} X^{\text{wt}(p)} t_{\varphi(p)}^{\frac{1}{2}} \left( \prod_{k \in f^{+}(p)} \frac{t_{\beta_{k}^{\vee}}^{-\frac{1}{2}} (1 - t_{\beta_{k}^{\vee}})}{1 - q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}} \right) \left( \prod_{k \in f^{-}(p)} \frac{t_{\beta_{k}^{\vee}}^{-\frac{1}{2}} (1 - t_{\beta_{k}^{\vee}}) q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}}{1 - q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}} \right),$$

where the sum is over the set  $\mathcal{B}(\vec{\mu}) = \mathcal{B}(1, \vec{w})$  of alcove walks of type  $i_1, \ldots, i_\ell$  beginning at 1.

Proof. Since  $E_{\mu} = \tau_{X^{\mu}m}^{\vee} \mathbf{1}$ ,

$$X^{\mathrm{end}(p)}\mathbf{1} = X^{\mathrm{wt}(p)}T_{\varphi(p)}^{\vee}\mathbf{1} = X^{\mathrm{wt}(p)}t_{\varphi(p)}^{\frac{1}{2}}\mathbf{1} \qquad \text{and} \qquad Y^{\lambda^{\vee}}\mathbf{1} = q^{\langle \lambda^{\vee}, \rho_c \rangle}\mathbf{1},$$

applying the formula for  $\tau_{X^{\mu}m}^{\vee}$  in Theorem 2.2 to 1 gives the formula in the statement.

Remark 3.2. From the expansion of  $E_{\mu}$  in Theorem 3.1, the nonsymmetric Macdonald polynomial  $E_{\mu}$  has top term  $t_m^{1/2}X^{\mu}$ , where  $X^{\mu}m$  is the minimal length representative of the coset  $X^{\mu}W_0$ . This term is the term corresponding to the unique alcove walk in  $\mathcal{B}(\vec{\mu})$  with no folds.

**Remark 3.3.** If  $w = X^{\mu}m = s_{i_1}^{\vee} \cdots s_{i_{\ell}}^{\vee}$  is a reduced word for the minimal length element of the coset  $X^{\mu}W_0$  then  $w^{-1} = s_{i_{\ell}}^{\vee} \cdots s_{i_1}^{\vee}$  is a walk from 1 to  $w^{-1}$  which stays completely in the dominant chamber. This has the effect that the roots  $\beta_{\ell}^{\vee}, \ldots, \beta_{1}^{\vee}$  are all of the form  $-\gamma^{\vee} + jd$  with  $\gamma^{\vee} \in (R^{\vee})^{+}$  (positive coroots) and  $j \in \mathbb{Z}_{>0}$ . The *height* of a coroot  $\gamma^{\vee}$  is

$$\operatorname{ht}(\gamma^{\vee}) = \langle \gamma^{\vee}, \rho \rangle, \quad \text{where } \rho = \frac{1}{2} \sum_{\alpha \in \mathbb{R}^+} \alpha.$$

In the case that all the parameters are equal  $(t_i = t = q^c \text{ for } i = 0, ..., n)$  the values which appear in Theorem 3.1,

$$q^{\langle -\beta_k^\vee, \rho_c \rangle} = q^{\langle \gamma^\vee - jd, \rho_c \rangle} = q^j t^{\operatorname{ht}(\gamma^\vee)}, \qquad \text{have positive exponents (in $\mathbb{Z}_{>0}$)}.$$

The set of dominant integral weights is

$$(\mathfrak{h}_{\mathbb{Z}}^*)^+ = \{ \mu \in \mathfrak{h}_{\mathbb{Z}}^* \mid \langle \mu, \alpha_i^{\vee} \rangle \ge 0 \text{ for } i = 1, \dots, n \}.$$
(3.8)

Recall the notation for  $t_w^{1/2}$  from (3.6). For  $\mu \in (\mathfrak{h}_{\mathbb{Z}}^*)^+$ , the symmetric Macdonald polynomial (see [Mac3, Remarks after (6.8)]) is

$$P_{\mu} = \mathbf{1}_{0} E_{\mu} \mathbf{1}$$
 where  $\mathbf{1}_{0} = \sum_{w \in W_{0}} t_{w_{0}w}^{-\frac{1}{2}} T_{w},$  (3.9)

so that  $T_i \mathbf{1}_0 = t_i^{1/2} \mathbf{1}_0$  for i = 1, 2 ... n, and  $\mathbf{1}_0$  has top term  $T_{w_0}$  with coefficient 1. The symmetric Macdonald polynomials are  $W_0$ -symmetric polynomials in  $X^{\mu}$  which are eigenvectors for the action of  $W_0$ -symmetric polynomials in the  $Y^{\lambda^{\vee}}$ .

**Theorem 3.4.** Let  $\mu \in (\mathfrak{h}_{\mathbb{Z}}^*)^+$  and let  $X^{\mu}m = s_{i_1}^{\vee} \cdots s_{i_\ell}^{\vee}$  be a reduced word for the minimal length element  $X^{\mu}m$  in the coset  $X^{\mu}W_0$ . Let  $\beta_{\ell}^{\vee}, \ldots, \beta_{1}^{\vee}$  be as defined in (2.35) and let

$$\mathcal{P}(\vec{\mu}) = \bigcup_{v \in W_0} \mathcal{B}(v, \vec{w})$$

be the set of alcove walks of type  $\vec{w} = (i_1, \dots, i_\ell)$  beginning at an element  $v \in W_0$ . Then the symmetric Macdonald polynomial

$$P_{\mu} = \sum_{p \in \mathcal{P}(\vec{\mu})} X^{\text{wt}(p)} t_{\varphi(p)}^{\frac{1}{2}} t_{w_{0}\iota(p)}^{-\frac{1}{2}} \left( \prod_{k \in f^{+}(p)} \frac{t_{\beta_{k}^{\vee}}^{-\frac{1}{2}} (1 - t_{\beta_{k}^{\vee}})}{1 - q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}} \right) \left( \prod_{k \in f^{-}(p)} \frac{t_{\beta_{k}^{\vee}}^{-\frac{1}{2}} (1 - t_{\beta_{k}^{\vee}}) q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}}{1 - q^{\langle -\beta_{k}^{\vee}, \rho_{c} \rangle}} \right),$$

where  $\iota(p)$  is the initial alcove of the path p.

*Proof.* The expression

$$\mathbf{1}_0 = \sum_{v \in W_0} t_{w_0 v}^{-\frac{1}{2}} X^v, \quad \text{gives} \quad P_{\mu} \mathbf{1} = \mathbf{1}_0 E_{\mu} \mathbf{1} = \sum_{v \in W_0} t_{w_0 v}^{-\frac{1}{2}} X^v \tau_{X^{\mu} m}^{\vee} \mathbf{1},$$

which is computed by the same method as in Theorem 2.2 and Theorem 3.1.

Remark 3.5. The Hall-Littlewood polynomials or Macdonald spherical functions are  $P_{\mu}(0,t)$  and the Schur functions or Weyl characters are  $s_{\mu} = P_{\mu}(0,0)$ . In the first case the formula in Theorem 3.4 reduces to the formula for the Macdonald spherical functions in terms of positively folded alcove walks as given in [Sc, Thm. 1.1] (see also [Ra, Thm. 4.2(a)]). In the case q = t = 0, the formula in Theorem 3.4 reduces to the formula for the Weyl characters in terms of maximal dimensional positively folded alcove walks (the Littelmann path model) as given in [GL, Cor. 1 p. 62], or the  $\lambda$ -chain formulation of [LP1],[LP2].

When  $\mu$  is not a regular weight, the formula for  $P_{\mu}$  in Theorem 3.4 has an alternate formulation as a sum over paths p whose initial alcove  $\iota(p)$  is in the minimal coset representatives  $W^{\mu}$  of  $W_0/W_{\mu}$ , where  $W_{\mu} = \{w \in W_0 \mid w\mu = \mu\}$  is the stabilzer of  $\mu$ . To see this, suppose  $s_i\mu = \mu$  for some  $i \in \{1, \ldots, n\}$ . Then  $\langle \mu, \alpha_i^{\vee} \rangle = 0$  implies  $Y^{-\alpha_i^{\vee}} E_{\mu} \mathbf{1} = t_i^{-1} E_{\mu} \mathbf{1}$ . Further, let  $X^{\mu}m$  be the minimal length element in the coset  $X^{\mu}W_0$ . Then  $s_iX^{\mu}m = X^{s_i\mu}s_im = X^{\mu}ms_j$  for some  $j \in \{1, \ldots, n\}$ , so  $\tau_i^{\vee} E_{\mu} \mathbf{1} = \tau_i^{\vee} \tau_{X^{\mu}m}^{\vee} \mathbf{1} = \tau_{X^{\mu}m}^{\vee} \tau_j^{\vee} \mathbf{1} = 0$ . Therefore,  $T_i E_{\mu} \mathbf{1} = \left(\tau_i^{\vee} - \frac{t_i^{-\frac{1}{2}} - t_i^{\frac{1}{2}}}{1 - Y^{-\alpha_i^{\vee}}}\right) E_{\mu} \mathbf{1} = t_i^{\frac{1}{2}} E_{\mu} \mathbf{1}$ , and

$$P_{\mu}\mathbf{1} = \mathbf{1}_{0}E_{\mu}\mathbf{1} = \sum_{w \in W_{0}} t_{w_{0}w}^{-\frac{1}{2}} T_{w} E_{\mu}\mathbf{1} = t_{w_{0}}^{\frac{1}{2}} \left(\sum_{v \in W^{\mu}} t_{w_{0}v}^{-\frac{1}{2}} T_{v}\right) \left(\sum_{u \in W_{\mu}} t_{w_{0}u}^{-\frac{1}{2}} T_{u}\right) E_{\mu}\mathbf{1}$$

$$= W_{\mu}(t) \sum_{v \in W^{\mu}} t_{w_{0}v}^{-\frac{1}{2}} X^{v} \tau_{X^{\mu}m}^{\vee} \mathbf{1},$$

where  $W_{\mu}(t) = \sum_{w \in W_{\mu}} t_w$  is the Poincaré polynomial of  $W_{\mu}$ .

**Remark 3.6.** In order to derive the Haglund-Haiman-Loehr formula for  $P_{\mu}$  in the type  $GL_n$  case, one would need to start with a similar formula that is based on the alcove walk from the fundamental alcove to the alcove  $X^{\mu}$  instead of using the formula in Theorem 3.4. Such a formula is equivalent to expanding the sum

$$P_{\mu}\mathbf{1} = t_{m}^{\frac{1}{2}} \sum_{v \in W_{0}} t_{w_{0}v}^{-\frac{1}{2}} X^{v} \tau_{X^{\mu}m}^{\vee} T_{m}^{-1} \mathbf{1}$$

in terms of the monomial basis.

Alternatively, let  $w_{\mu} = s_{i_1} \cdots s_{i_l}$  be the longest element in  $W_{\mu}$ , so that  $w_0 w_{\mu}$  is the shortest element in the coset  $w_0 W_{\mu}$ . Renormalizing  $P_{\mu}$ ,

$$\left(\prod_{j=1}^{l} t_{i_{j}}^{\frac{1}{2}} + \frac{t_{i_{j}}^{-\frac{1}{2}} - t_{i_{j}}^{\frac{1}{2}}}{1 - q^{-\langle \alpha_{i_{j}}^{\vee}, X^{-\mu} m \cdot \rho_{c} \rangle}}\right) P_{\mu} \mathbf{1} = \mathbf{1}_{0} \left(\prod_{j=1}^{l} t_{i_{j}}^{\frac{1}{2}} + \frac{t_{i_{j}}^{-\frac{1}{2}} - t_{i_{j}}^{\frac{1}{2}}}{1 - q^{-\langle \alpha_{i_{j}}^{\vee}, X^{-\mu} m \cdot \rho_{c} \rangle}}\right) E_{\mu} \mathbf{1}$$

$$= \mathbf{1}_{0} \prod_{j=1}^{l} \left(T_{i_{j}}^{\frac{1}{2}} + \frac{t_{i_{j}}^{-\frac{1}{2}} - t_{i_{j}}^{\frac{1}{2}}}{1 - Y^{-\alpha_{i_{j}}^{\vee}}}\right) E_{\mu} \mathbf{1} = \mathbf{1}_{0} \tau_{w_{0}w_{\mu}}^{\vee} E_{\mu} \mathbf{1} = \mathbf{1}_{0} E_{w_{0}w_{\mu}\mu} \mathbf{1},$$

and expanding the right hand side in the same way as in the proof of Theorem 3.4 will produce a formula based on the alcove walk from the fundamental alcove to the alcove  $X^{w_0w_\mu\mu}$ . This walk has the same length as the walk from the fundamental alcove to the alcove  $X^\mu$ .

### 4 Examples

#### 4.1 Type $A_1$

The Weyl group  $W_0 = \langle s_1 \mid s_1^2 = 1 \rangle$  has order two and acts on the lattices

$$\mathfrak{h}_{\mathbb{Z}} = \mathbb{Z}\omega^{\vee} \quad \text{and} \quad \mathfrak{h}_{\mathbb{Z}}^* = \mathbb{Z}\omega \quad \text{by} \quad s_1\omega^{\vee} = -\omega^{\vee} \quad \text{and} \quad s_1\omega = -\omega,$$
 (4.1)

and

$$\varphi^{\vee} = \alpha^{\vee} = 2\omega^{\vee}, \qquad \varphi = \alpha = 2\omega, \quad \text{and} \quad \langle \omega^{\vee}, \alpha \rangle = 1.$$
 (4.2)

The double affine braid group  $\widetilde{\mathcal{B}}$  is generated by  $T_0, T_1, g, X^{\omega}$ , and  $q^{1/2}$ , with relations

$$T_0 = gT_1g^{-1}, \qquad g^2 = 1, \qquad q = X^{\delta},$$
 
$$qX^{\omega} = q^{1/2}X^{-\omega}q, \qquad T_1X^{\omega}T_1 = X^{-\omega}, \quad \text{and} \quad T_0X^{-\omega}T_0 = q^{-1}X^{\omega}.$$
 (4.3)

In the double affine braid group

$$g = Y^{\omega^{\vee}} T_1^{-1}, \qquad T_0 = Y^{\varphi^{\vee}} T_1^{-1}, \qquad g^{\vee} = X^{\omega} T_1, \qquad (T_0^{\vee})^{-1} = X^{\varphi} T_1^{\vee}.$$
 (4.4)

At this point, the following Proposition, which is the Type  $A_1$  case of Theorem 2.1, is easily proved by direct computation.

**Proposition 4.1.** (Duality). Let  $Y^d = q^{-1}$ . The double affine braid group  $\widetilde{\mathcal{B}}$  is generated by  $T_0^{\vee}, T_1^{\vee}, g^{\vee}, Y^{\omega^{\vee}}$  and  $q^{1/2}$  with relations

$$Y^d = q^{-1}, \qquad (g^{\vee})^2 = 1, \qquad T_0^{\vee} = g^{\vee} T_1^{\vee} (g^{\vee})^{-1},$$
 
$$g^{\vee} Y^{\omega^{\vee}} = q^{-1/2} Y^{-\omega^{\vee}} g^{\vee}, \quad T_1^{-1} Y^{\omega^{\vee}} T_1^{-1} = Y^{-\omega^{\vee}}, \quad and \quad (T_0^{\vee})^{-1} Y^{-\omega^{\vee}} (T_0^{\vee})^{-1} = q Y^{\omega^{\vee}}.$$

The double affine Hecke algebra  $\widetilde{H}$  is  $\mathbb{C}\widetilde{\mathcal{B}}$  with the additional relations

$$T_i^2 = (t^{1/2} - t^{-1/2})T_i + 1$$
, for  $i = 0, 1$ , and  $t_0 = t_1 = t = q^c$ . (4.5)

Using (4.5), the relations in Proposition (4.1) give

$$g^{\vee}Y^{\omega^{\vee}} = q^{-1/2}Y^{-\omega^{\vee}}g^{\vee}, \qquad T_{1}Y^{\omega^{\vee}} = Y^{-\omega^{\vee}}T_{1} + (t^{1/2} - t^{-1/2})\frac{Y^{\omega^{\vee}} - Y^{-\omega^{\vee}}}{1 - Y^{-\alpha^{\vee}}}, \quad \text{and}$$
$$T_{0}^{\vee}Y^{\omega^{\vee}} = q^{-1}Y^{-\omega^{\vee}}T_{0}^{\vee} + (t^{1/2} - t^{-1/2})\left(\frac{Y^{\omega^{\vee}} - q^{-1}Y^{-\omega^{\vee}}}{1 - qY^{\alpha^{\vee}}}\right).$$

With  $Y^{-\alpha_0^{\vee}} = qY^{\alpha^{\vee}}$  and  $Y^{-\alpha_1^{\vee}} = Y^{-\alpha^{\vee}}$ , then

$$\tau_g^{\vee} = g^{\vee}, \quad \text{and} \quad \tau_i^{\vee} = T_i^{\vee} - (t^{1/2} - t^{-1/2}) \left(\frac{1}{1 - Y^{-\alpha_i^{\vee}}}\right), \quad \text{for } i = 0, 1.$$

To illustrate Theorem 2.2, note that  $X^{-2\omega} = s_1^{\vee} s_0^{\vee}$  is a reduced word and

$$\begin{split} \tau_1^\vee \tau_0^\vee &= \left(T_1^\vee + \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_1^\vee}}\right)\tau_0^\vee \\ &= T_1^\vee T_0^\vee + T_1^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_0^\vee}} + (T_0^\vee)^{-1}\frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^\vee}} + \left(\frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^\vee}}\right) \left(\frac{t^{-1/2}(1-t)Y^{-\alpha_0^\vee}}{1-Y^{-\alpha_0^\vee}}\right) \\ &= X^{-2\omega} + T_1^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_0^\vee}} + X^{2\omega}T_1^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^\vee}} + \left(\frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^\vee}}\right) \left(\frac{t^{-1/2}(1-t)Y^{-\alpha_0^\vee}}{1-Y^{-\alpha_0^\vee}}\right). \end{split}$$

The corresponding paths in  $\mathcal{B}(1, \overrightarrow{-2\omega}) = \mathcal{B}(\overrightarrow{-2\omega})$  are

$$X^{-2\omega} \qquad T_1^{\vee} \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_0^{\vee}}} \qquad X^{2\omega} T_1^{\vee} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^{\vee}}} \qquad \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_1^{\vee}}} \frac{t^{-1/2}(1-t)Y^{-\alpha_0^{\vee}}}{1-Y^{-s_0\alpha_1^{\vee}}}$$

The polynomial representation is defined by

$$T_i \mathbf{1} = t^{1/2} \mathbf{1}$$
, and  $g \mathbf{1} = \mathbf{1}$ .

In this case

$$\rho_c = \frac{1}{2}c\alpha \quad \text{and} \quad W^0 = \{X^{-\ell\omega} \mid \ell \in \mathbb{Z}_{\geq 0}\} \cup \{X^{\ell\omega}s_1^{\vee} \mid \ell \in \mathbb{Z}_{> 0}\}, \tag{4.6}$$

is the set of minimal length coset representatives of  $W^{\vee}/W_0$ .

Applying the expansion of  $\tau_1^{\vee}\tau_0^{\vee}$  to **1** and using

$$Y^{-\alpha_0^\vee}\mathbf{1} = qY^{\alpha^\vee}\mathbf{1} = qq^c\mathbf{1} = qt\mathbf{1}, \qquad \text{and} \qquad Y^{-s_0\alpha_1^\vee}\mathbf{1} = Y^{\alpha^\vee+2d}\mathbf{1} = q^2Y^{\alpha^\vee}\mathbf{1} = q^2t,$$

gives

$$\begin{split} E_{-2\omega} &= \tau_1^{\vee} \tau_0^{\vee} \mathbf{1} \\ &= X^{-2\omega} + t^{1/2} \frac{t^{-1/2} (1-t)}{1-qt} + X^{2\omega} t^{1/2} \frac{t^{-1/2} (1-t)}{1-q^2 t} + \left( \frac{t^{-1/2} (1-t)}{1-q^2 t} \right) \left( \frac{t^{-1/2} (1-t)qt}{1-qt} \right) \\ &= X^{-2\omega} + \frac{1-t}{1-qt} + X^{2\omega} \frac{1-t}{1-q^2 t} + \left( \frac{1-t}{1-q^2 t} \right) \left( \frac{(1-t)q}{1-qt} \right). \end{split}$$

Since  $\mathbf{1}_0 = T_1^{\vee} + t^{-1/2}$  the symmetric Macdonald polynomial  $P_{2\omega} = \mathbf{1}_0 E_{2\omega} = \mathbf{1}_0 \tau_0^{\vee} \mathbf{1}$  is

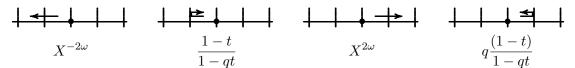
$$P_{2\omega} = \mathbf{1}_{0} E_{2\omega} = (T_{1}^{\vee} + t^{-1/2}) \tau_{0}^{\vee} \mathbf{1}$$

$$= \left( T_{1}^{\vee} T_{0}^{\vee} + T_{1}^{\vee} \frac{t^{-1/2} (1 - t)}{1 - Y^{-\alpha_{0}^{\vee}}} + t^{-1/2} (T_{0}^{\vee})^{-1} + t^{-1/2} \frac{t^{-1/2} (1 - t) Y^{-\alpha_{0}^{\vee}}}{1 - Y^{-\alpha_{0}^{\vee}}} \right) \mathbf{1}$$

$$= \left( X^{-2\omega} + t^{1/2} \frac{t^{-1/2} (1 - t)}{1 - qt} + t^{-1/2} X^{2\omega} T_{1}^{\vee} + t^{-1/2} \frac{t^{-1/2} (1 - t) qt}{1 - qt} \right) \mathbf{1}$$

$$= \left( X^{-2\omega} + \frac{1 - t}{1 - qt} + X^{2\omega} + \frac{(1 - t)q}{1 - qt} \right) \mathbf{1} = \left( X^{2\omega} + X^{-2\omega} + (1 + q) \frac{1 - t}{1 - qt} \right) \mathbf{1}.$$

The corresponding paths in  $\mathcal{P}(\overrightarrow{2\omega})$  are



### 4.2 Type $A_2$

The Weyl group  $W_0 = \langle s_1, s_2 | s_1^2 = s_2^2 = 1, s_1 s_2 s_1 = s_2 s_1 s_2 \rangle$  acts on the lattices

$$\mathfrak{h}_{\mathbb{Z}} = \mathbb{Z}\omega_1^{\vee} + \mathbb{Z}\omega_2^{\vee} \quad \text{and} \quad \mathfrak{h}_{\mathbb{Z}}^* = \mathbb{Z}\omega_1 + \mathbb{Z}\omega_2,$$
(4.7)

where  $s_1$  and  $s_2$  are the reflections in the hyperplanes determined by

$$\alpha_1^{\vee} = 2\omega_1^{\vee} - \omega_2^{\vee}, \quad \alpha_2^{\vee} = -\omega_1^{\vee} + 2\omega_2^{\vee}, \quad \alpha_1 = 2\omega_1 - \omega_2, \quad \text{and} \quad \alpha_2 = -\omega_1 + 2\omega_2,$$
 (4.8)

with  $\langle \omega_i^{\vee}, \alpha_j \rangle = \delta_{ij}$ , and  $\langle \omega_i, \alpha_j^{\vee} \rangle = \delta_{ij}$ . In this case,

$$\varphi^{\vee} = \alpha_1^{\vee} + \alpha_2^{\vee}, \quad \text{and} \quad \varphi = \alpha_1 + \alpha_2.$$
 (4.9)

The double affine braid group  $\widetilde{\mathcal{B}}$  is generated by  $T_0, T_1, T_2, g, X^{\omega_1}, X^{\omega_2}$ , and  $q^{1/3}$ , with relations

$$T_{i}T_{j}T_{i} = T_{j}T_{i}T_{j}, \quad \text{for } i \neq j,$$

$$X^{\mu}X^{\lambda} = X^{\mu+\lambda} = X^{\lambda}X^{\mu}, \quad \text{for } \mu, \lambda \in \mathfrak{h}_{\mathbb{Z}}^{*},$$

$$T_{1}X^{\omega_{2}} = X^{\omega_{2}}T_{1}, \quad T_{2}X^{\omega_{1}} = X^{\omega_{1}}T_{2}, \quad T_{1}X^{\omega_{1}}T_{1} = X^{-\omega_{1}+\omega_{2}}, \quad T_{2}X^{\omega_{2}}T_{2} = X^{\omega_{1}-\omega_{2}}, \quad (4.10)$$

$$g^{3} = 1, \quad gX^{\omega_{1}} = q^{1/3}X^{-\omega_{1}+\omega_{2}}g, \quad gX^{\omega_{2}} = q^{2/3}X^{-\omega_{1}}g,$$

$$gT_{0}g^{-1} = T_{1}, \quad gT_{1}g^{-1} = T_{2}, \quad gT_{2}g^{-1} = T_{0}.$$

The formula (2.27) gives

$$g = Y^{\omega_1^{\vee}} T_1^{-1} T_2^{-1}, \qquad g^2 = Y^{\omega_2^{\vee}} T_2^{-1} T_1^{-1}, \qquad T_0 = Y^{\varphi^{\vee}} T_1^{-1} T_2^{-1} T_1^{-1}, \tag{4.11}$$

$$g^{\vee} = X^{\omega_1} T_1^{\vee} T_2^{\vee}, \qquad (g^{\vee})^2 = X^{\omega_2} T_2^{\vee} T_1^{\vee}, \qquad (T_0^{\vee})^{-1} = X^{\varphi} T_1^{\vee} T_2^{\vee} T_1^{\vee}. \tag{4.12}$$

At this point, the following Proposition, which is the Type  $A_2$  case of Theorem 2.1, is easily proved by direct computation.

**Proposition 4.2.** (Duality). Let  $Y^d = q^{-1}$ . The double affine braid group  $\widetilde{\mathcal{B}}$  is generated by  $T_0^{\vee}$ ,  $T_1^{\vee}$ ,  $T_2^{\vee}$ ,  $g^{\vee}$ ,  $Y^{\omega_1^{\vee}}$ ,  $Y^{\omega_2^{\vee}}$  and  $q^{1/3}$ , with relations

$$\begin{split} (T_1^\vee)^{-1}Y^{\omega_1^\vee}(T_1^\vee)^{-1} &= Y^{-\omega_1^\vee+\omega_2^\vee}, \quad (T_2^\vee)^{-1}Y^{\omega_2^\vee}(T_2^\vee)^{-1} &= Y^{\omega_1^\vee-\omega_2^\vee}, \\ (T_1^\vee)^{-1}Y^{\omega_2^\vee} &= Y^{\omega_2^\vee}(T_1^\vee)^{-1}, \quad (T_2^\vee)^{-1}Y^{\omega_1^\vee} &= Y^{\omega_1^\vee}(T_2^\vee)^{-1}, \\ (g^\vee)^3 &= 1, \quad g^\vee Y^{\omega_1^\vee} &= q^{-1/3}Y^{-\omega_1^\vee+\omega_2^\vee}g^\vee, \quad g^\vee Y^{\omega_2^\vee} &= q^{-2/3}Y^{-\omega_1^\vee}g^\vee, \\ g^\vee T_0^\vee(g^\vee)^{-1} &= T_1^\vee, \qquad g^\vee T_1^\vee(g^\vee)^{-1} &= T_2^\vee \quad \text{ and } \quad g^\vee T_2^\vee(g^\vee)^{-1} &= T_0^\vee. \end{split}$$

To give a concrete example of Theorem 3.4 let us compute the symmetric Macdonald polynomial  $P_{\rho}$  where  $\rho = \alpha_1 + \alpha_2$ . Since

$$\mathbf{1_0} = X^{s_1 s_2 s_1} + t^{-1/2} X^{s_1 s_2} + t^{-1/2} X^{s_2 s_1} + t^{-2/2} X^{s_1} + t^{-2/2} X^{s_2} + t^{-3/2},$$

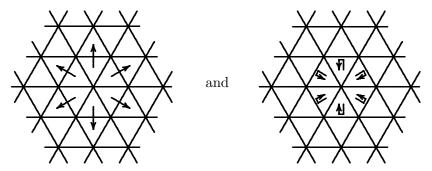
and  $X^{\rho}m = s_0^{\vee}$  is the minimal length element of the coset  $X^{\rho}W_0$ ,

$$\begin{split} P_{\rho} &= \mathbf{1}_{0} E_{\rho} = \mathbf{1}_{0} \tau_{0}^{\vee} \mathbf{1} \\ &= \left( X^{s_{1}s_{2}s_{1}} + t^{-1/2} X^{s_{1}s_{2}} + t^{-1/2} X^{s_{2}s_{1}} \right) \left( T_{0}^{\vee} + \frac{t^{-1/2} (1-t)}{1-Y^{-\alpha_{0}^{\vee}}} \right) \mathbf{1} \\ &+ \left( t^{-2/2} X^{s_{1}} + t^{-2/2} X^{s_{2}} + t^{-3/2} \right) \left( (T_{0}^{\vee})^{-1} + \frac{t^{-1/2} (1-t) Y^{-\alpha_{0}^{\vee}}}{1-Y^{-\alpha_{0}^{\vee}}} \right) \mathbf{1} \\ &= \left( X^{s_{1}s_{2}s_{1}s_{0}} + t^{-1/2} X^{s_{1}s_{2}s_{0}} + t^{-1/2} X^{s_{2}s_{1}s_{0}} + t^{-2/2} X^{s_{1}s_{0}} + t^{-2/2} X^{s_{2}s_{0}} + t^{-3/2} X^{s_{0}} \right) \mathbf{1} \\ &+ \left( X^{s_{1}s_{2}s_{1}} + t^{-1/2} X^{s_{1}s_{2}} + t^{-1/2} X^{s_{2}s_{1}} \right) \frac{t^{-1/2} (1-t)}{1-Y^{-\alpha_{0}^{\vee}}} \mathbf{1} \\ &+ \left( t^{-2/2} X^{s_{1}} + t^{-2/2} X^{s_{2}} + t^{-3/2} \right) \frac{t^{-1/2} (1-t) Y^{-\alpha_{0}^{\vee}}}{1-Y^{-\alpha_{0}^{\vee}}} \mathbf{1}. \end{split}$$

Since 
$$Y^{-\alpha_0^{\vee}} \mathbf{1} = Y^{\varphi^{\vee} - d} \mathbf{1} = q Y^{\alpha_1^{\vee} + \alpha_2^{\vee}} \mathbf{1} = q t^2 \mathbf{1}$$
,

$$\begin{split} P_{\rho} &= \left( \begin{array}{c} X^{w_0\rho} + t^{-1/2} X^{s_1s_2\rho} T_2^{\vee} + t^{-1/2} X^{s_2s_1\rho} T_1^{\vee} \\ + t^{-2/2} X^{s_1\rho} T_2^{\vee} T_1^{\vee} + t^{-2/2} X^{s_2\rho} T_1^{\vee} T_2^{\vee} + t^{-3/2} X^{\rho} T_1^{\vee} T_2^{\vee} T_1^{\vee} \end{array} \right) \mathbf{1} \\ &+ \left( T_1^{\vee} T_2^{\vee} T_1^{\vee} + t^{-1/2} T_1^{\vee} T_2^{\vee} + t^{-1/2} T_2^{\vee} T_1^{\vee} \right) \frac{t^{-1/2} (1-t)}{1-qt^2} \mathbf{1} \\ &+ \left( t^{-2/2} T_1^{\vee} + t^{-2/2} T_2^{\vee} + t^{-3/2} \right) \frac{t^{-1/2} (1-t)qt^2}{1-qt^2} \mathbf{1} \\ &= \left( X^{w_0\rho} + X^{s_1s_2\rho} + X^{s_2s_1\rho} + X^{s_1\rho} + X^{s_2\rho} + X^{\rho} \right) \mathbf{1} \\ &+ \left( t^{3/2} + t^{1/2} + t^{1/2} \right) \frac{t^{-1/2} (1-t)}{1-qt^2} \mathbf{1} + \left( t^{-1/2} + t^{-1/2} + t^{-3/2} \right) \frac{t^{-1/2} (1-t)qt^2}{1-qt^2} \mathbf{1} \\ &= \left( X^{w_0\rho} + X^{s_1s_2\rho} + X^{s_2s_1\rho} + X^{s_1\rho} + X^{s_2\rho} + X^{\rho} + (t+2+2qt+q) \frac{1-t}{1-qt^2} \right) \mathbf{1}. \end{split}$$

The set  $\mathcal{P}(\vec{\rho})$  contains 12 alcove walks,



The Hall-Littlewood polynomial and the Weyl character are

$$P_{\rho}(0,t) = m_{\rho} + (2+t)(1-t)$$
 and  $s_{\rho} = P_{\rho}(0,0) = m_{\rho} + 2$ ,

where  $m_{\rho} = X^{w_0\rho} + X^{s_1s_2\rho} + X^{s_2s_1\rho} + X^{s_1\rho} + X^{s_2\rho} + X^{\rho}$ .

The expression  $X^{s_1s_2\rho}s_2 = s_1^{\vee} s_2^{\vee} s_0^{\vee}$  is a reduced word for the minimal length element in the coset  $X^{s_1s_2\rho}W_0$  and Theorem 2.2 is illustrated by

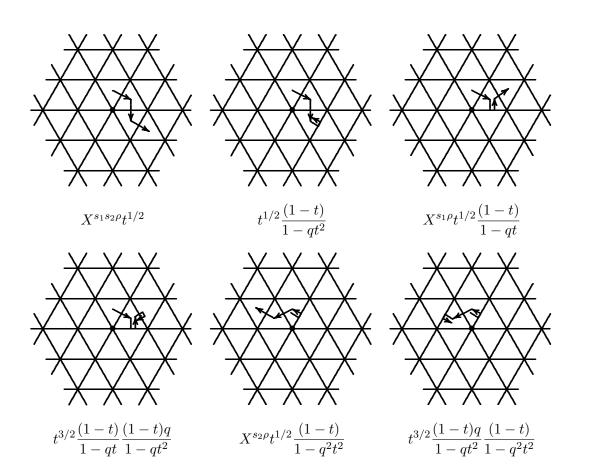
$$\begin{split} \tau_1^\vee \tau_2^\vee \tau_0^\vee &= \left(T_1^\vee + \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_1^\vee}}\right) \tau_1^\vee \tau_0^\vee = \left(T_1^\vee \tau_2^\vee + \tau_2^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_2\alpha_1^\vee}}\right) \tau_0^\vee \\ &= \left(T_1^\vee T_2^\vee + T_1^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_2^\vee}} + T_2^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_2\alpha_1^\vee}} + \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_2^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_2\alpha_1^\vee}}\right) \tau_0^\vee \\ &= T_1^\vee T_2^\vee \tau_0^\vee + T_1^\vee \tau_0^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_2^\vee}} + T_2^\vee \tau_0^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} + \tau_0^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \\ &= T_1^\vee T_2^\vee T_0^\vee + T_1^\vee T_2^\vee \frac{t^{-1/2}(1-t)}{1-Y^{-\alpha_0^\vee}} + T_1^\vee (T_0^\vee)^{-1} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_2^\vee}} + T_2^\vee (T_0^\vee)^{-1} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \\ &+ T_1^\vee \frac{t^{-1/2}(1-t)Y^{-\alpha_0^\vee}}{1-Y^{-s_0}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} + (T_0^\vee)^{-1} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \\ &+ T_2^\vee \frac{t^{-1/2}(1-t)Y^{-\alpha_0^\vee}}{1-Y^{-s_0s_2\alpha_1^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} + (T_0^\vee)^{-1} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0\alpha_2^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \\ &+ \frac{t^{-1/2}(1-t)Y^{-\alpha_0^\vee}}{1-Y^{-s_0s_2\alpha_1^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}} \frac{t^{-1/2}(1-t)}{1-Y^{-s_0s_2\alpha_1^\vee}}, \end{split}$$

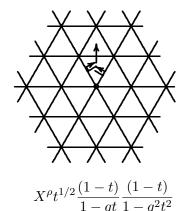
where the eight terms in this expansion correspond to the eight alcove walks in  $\mathcal{B}(1, s_1^{\vee} s_2^{\vee} s_0^{\vee}) = \mathcal{B}(\overrightarrow{s_1 s_2 \rho})$  pictured below. Applying the expansion of  $\tau_1^{\vee} \tau_2^{\vee} \tau_0^{\vee}$  to **1** and using

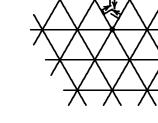
$$Y^{-\alpha_0^{\vee}} \mathbf{1} = Y^{\varphi^{\vee} - d} \mathbf{1} = qt^2 \mathbf{1}, \qquad Y^{-s_0 \alpha_2^{\vee}} \mathbf{1} = Y^{\alpha_1^{\vee} - d} \mathbf{1} = qt \mathbf{1},$$
and
$$Y^{-s_0 s_2 \alpha_1^{\vee}} \mathbf{1} = Y^{\varphi^{\vee} - 2d} \mathbf{1} = q^2 t^2 \mathbf{1},$$
(4.13)

computes

$$\begin{split} E_{s_1s_2\rho} &= \left( X^{s_1s_2\rho} t^{1/2} + t \frac{t^{-1/2}(1-t)}{1-qt^2} + X^{s_1\rho} t \frac{t^{-1/2}(1-t)}{1-qt} + t^{1/2} \frac{t^{-1/2}(1-t)qt^2}{1-qt^2} \frac{t^{-1/2}(1-t)}{1-qt} \right. \\ &\quad + X^{s_2\rho} t \frac{t^{-1/2}(1-t)}{1-q^2t^2} + t^{1/2} \frac{t^{-1/2}(1-t)qt^2}{1-qt^2} \frac{t^{-1/2}(1-t)}{1-q^2t^2} \\ &\quad + X^{\rho} t^{3/2} \frac{t^{-1/2}(1-t)}{1-qt} \frac{t^{-1/2}(1-t)}{1-qt} \frac{t^{-1/2}(1-t)qt^2}{1-qt^2} \frac{t^{-1/2}(1-t)}{1-qt} \frac{t^{-1/2}(1-t)}{1-qt} \frac{t^{-1/2}(1-t)}{1-qt^2} \right) \mathbf{1}. \\ &= t^{1/2} \left( X^{s_1s_2\rho} + \frac{(1-t)}{1-qt^2} + X^{s_1\rho} \frac{(1-t)}{1-qt} + t \frac{(1-t)q}{1-qt^2} \frac{(1-t)}{1-qt} + X^{s_2\rho} \frac{(1-t)}{1-qt^2} \right. \\ &\quad + t \frac{(1-t)q}{1-qt^2} \frac{(1-t)}{1-q^2t^2} + X^{\rho} \frac{(1-t)}{1-qt} \frac{(1-t)}{1-qt^2} + \frac{(1-t)q}{1-qt^2} \frac{(1-t)}{1-qt} \frac{(1-t)}{1-qt^2} \right) \mathbf{1}. \end{split}$$







$$t^{3/2} \frac{(1-t)}{1-qt} \frac{(1-t)q}{1-qt^2} \frac{(1-t)}{1-q^2t^2}$$

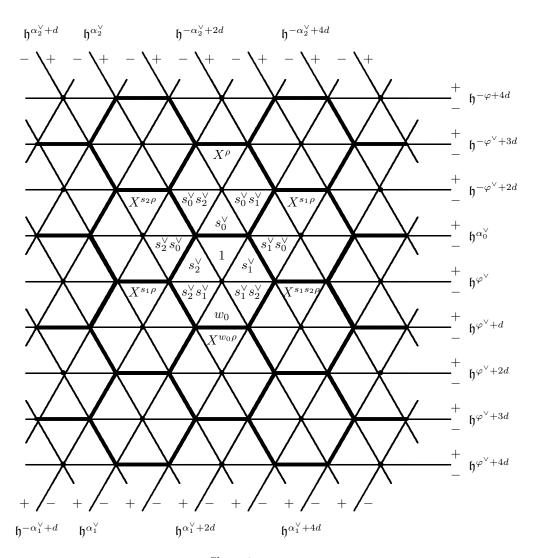
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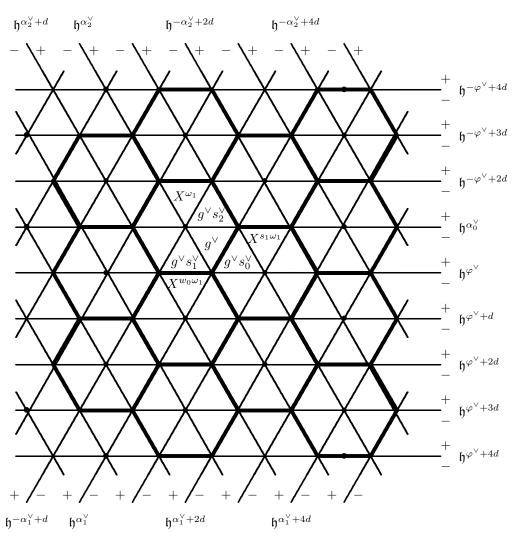
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# 5 Appendix: The bijection between W and alcoves in type $SL_3$

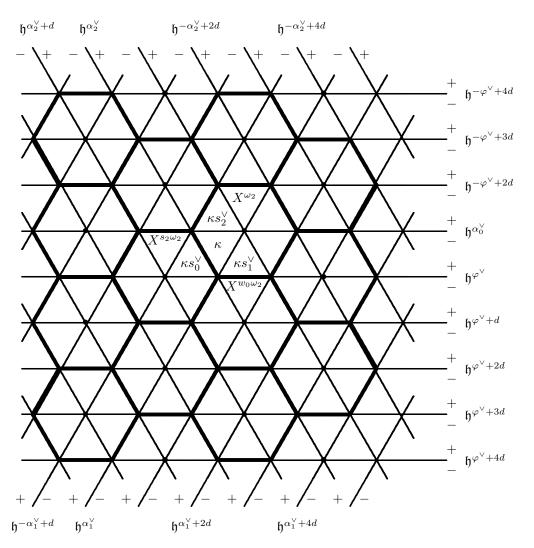
The following pictures illustrate the bijection of (2.19) for type  $SL_3$ . In this case,  $\Omega^{\vee} = \{1, g^{\vee}, (g^{\vee})^2\} \cong \mathbb{Z}/3\mathbb{Z}$ , and  $\Omega^{\vee} \times \mathfrak{h}_{\mathbb{R}}^*$  has 3 sheets. The alcoves are the triangles and the (centres of) hexagons are the elements of  $\mathfrak{h}_{\mathbb{Z}}^*$ .



Sheet 1



Sheet  $g^{\vee}$ 



Sheet  $\kappa = (g^{\vee})^2$